

Brillouin and Langevin functions

Langevin model

Model aim: The Langevin model describes the energy distribution of electrons of different spin direction (the spin distribution) in an external magnetic field

If θ is the angle between magnetic moment μ at H , the magnetic energy will be

$$E_m = -\mu \cdot H \cdot \cos(\theta) \quad (1)$$

The number of electrons of energy E_m follows the Boltzmann statistics

$$n(E_m) \sim e^{\frac{E_m}{kT}} \quad (2)$$

Number of electrons which magnetic moment between θ and $\theta + d\theta$, ϕ and $\phi + d\phi$,

$$dn(\theta, \phi) = A \cdot e^{\frac{-\mu \cdot H \cdot \cos(\theta)}{kT}} \cdot d\theta \cdot d\phi \quad (3)$$

Integration over ϕ gives

$$dn(\theta) = A \cdot e^{\frac{-\mu \cdot H \cdot \cos(\theta)}{kT}} \cdot 2\pi \cdot \sin(\theta) d\theta \quad (4)$$

Eq. 4 is normalized to the total number n_d of d- electrons

$$n_d = 2\pi \cdot A \cdot \int_{-\pi}^{\pi} e^{\frac{-\mu \cdot H \cdot \cos(\theta)}{kT}} \sin(\theta) d\theta = 2\pi \cdot A \cdot \int_{-1}^1 e^{\frac{-\mu \cdot H \cdot \cos(\theta)}{kT}} d(\cos(\theta)) \quad (5)$$

Integration gives

$$n_d = \frac{4\pi \cdot A \cdot kT}{\mu \cdot H} \sinh\left(\frac{\mu \cdot H}{kT}\right) \quad (6)$$

Therefore, distribution of electrons over angles is

$$dn(\theta) = \frac{4\pi \cdot n_d \cdot kT}{\mu \cdot H} \sinh\left(\frac{\mu \cdot H}{kT}\right) \cdot e^{\frac{-\mu \cdot H \cdot \cos(\theta)}{kT}} \cdot 2\pi \cdot \sin(\theta) d\theta \quad (7)$$

The magnetic moments perpendicular to H will cancel on the average, and the total magnetic moment m along H is calculated as

$$m = \int_{-\pi}^{\pi} \mu \cdot \cos(\theta) dn(\theta) = A \cdot 2\pi \int_{-\pi}^{\pi} \mu \cdot \cos(\theta) e^{\frac{-\mu \cdot H \cdot \cos(\theta)}{kT}} \cdot \sin(\theta) d\theta \quad (8)$$

The integration by parts gives

$$\frac{M}{M_0} = L\left(\frac{\mu \cdot H}{kT}\right) \quad (9)$$

where M is the saturation magnetization and L(x) is the Langevin function

$$L(x) = \coth(x) - \frac{1}{x} \quad (11)$$

$$\coth(x) = \frac{e^x + e^{-x}}{e^x - e^{-x}} \quad (12)$$

Ferromagnetism

Mean field. Weiss theory.

$$F = H + I \quad (20)$$

where F is the total magnetic field, I is the exchange field due to exchange interaction and the internal magnetic field, which is induced by the magnetic moments.

The exchange field I is changed proportionally to the change of magnetization

$$\frac{I}{I_0} = \frac{M}{M_0} \quad (21)$$

where I₀ is the exchange field at saturated magnetic moment.

Substitution of Eqs.9,11 into Eq.21 gives

$$\frac{I}{I_0} = L\left(\frac{\mu \cdot F}{kT}\right) = \coth\left(\frac{\mu \cdot F}{kT}\right) - \frac{kT}{\mu \cdot F} \quad (22)$$

Substitution of Eqs.20 into Eq.22 gives

$$\frac{I}{I_0} = L\left(\frac{\mu \cdot (H + I)}{kT}\right) \quad (23)$$

Or without external field

$$\frac{I}{I_0} = L\left(\frac{\mu \cdot I}{kT}\right) \quad (24)$$

Curie temperature

It is the case when the magnetization becomes zero $M=0$ and therefore the exchange field becomes zero as well $I=0$

Since

$$\lim_{x \rightarrow 0} (L(x)) = \frac{1}{3}x \quad (25)$$

Eq. 24 can be transformed at $I \rightarrow 0$

$$\lim_{x \rightarrow 0} [L(x)] = \frac{e^x + e^{-x}}{e^x - e^{-x}} - \frac{1}{x} = \frac{2}{2x + \frac{x^3}{3}} - \frac{1}{x} = \frac{1}{x} \left(\frac{2}{2 + \frac{x^2}{3}} - 1 \right) = -\frac{1}{x} \frac{x^2/3}{2+x} = -\frac{x}{6} \quad (11)$$

Using Eq.24 at small I gives

$$\frac{I}{I_0} = \frac{1}{3} \frac{\mu \cdot I}{kT_c} \quad (26)$$

Where T_c is Curie temperature

$$T_c = \frac{\mu \cdot I_0}{3k} \quad (26)$$

As I is linearly proportional to M

$$I = I_0 \frac{M}{M_0} = \frac{M}{M_0} \frac{3k \cdot T_c}{\mu} \quad (24)$$

$$\frac{M}{M_0} = L \left(\frac{\mu}{kT} \frac{M}{M_0} \frac{3k \cdot T_c}{\mu} \right) = L \left(3 \frac{M}{M_0} \frac{T_c}{T} \right) \quad (24)$$

With external magnetic field

$$\frac{M}{M_0} = L \left(\frac{\mu \cdot (I + H)}{kT} \right) = L \left(\frac{\mu \cdot \left(\frac{M}{M_0} \frac{3k \cdot T_c}{\mu} + H \right)}{kT} \right) = L \left(3 \frac{M}{M_0} \frac{T_c}{T} + H \frac{\mu}{kT} \right) \quad (23)$$

Derivative

$$\frac{\partial}{\partial H} \left(\frac{M}{M_0} \right) = \frac{\partial L}{\partial x} (x_0) \left[3 \frac{T_c}{T} \frac{\partial}{\partial H} \left(\frac{M}{M_0} \right) + \frac{\mu}{kT} \right] \quad (23)$$

$$x_0 = 3 \frac{M}{M_0} \frac{T_c}{T} + H \frac{\mu}{kT}$$

$$\frac{1}{M_0} \frac{\partial M}{\partial H} \left(1 - 3 \frac{T_c}{T} \frac{\partial L}{\partial x} \right) = \frac{\mu}{kT} \frac{\partial L}{\partial x} \quad (23)$$